

# Branes at toric conical singularities

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# Motivation

- We wish to understand the  $AdS/CFT$  correspondence away from the maximally supersymmetric case ( $AdS_5 \times S^5$ ), but in situations where dual theories can still be under control ( $\mathcal{N} = 1$ )

## Plan:

- Discuss the relevant supergravity backgrounds
- Elaborate on certain geometric aspects of the problem (new resolution parameters)

# Supergravity backgrounds and interpretation

# Supergravity solutions. I.

- We start with type IIB theory defined on  $\mathbb{R}^{3,1} \times Y$ , where  $Y$  is a Calabi-Yau threefold
- We place D3-branes with worldvolume  $\mathbb{R}^{3,1}$  at an isolated singular point of  $Y$
- We will consider singularities that have the form of complex cone over a surface  $X$
- We will assume that  $X$  is toric (hence  $Y$  is toric), which means that  $Y$  has at least a  $U(1)^3$  worth of isometries

# Supergravity solutions. II.

- Within supergravity these branes lead to solutions of the form

$$ds^2 = h^{-1/2}(y) \sum_{\mu=0}^3 dx_{\mu} dx_{\mu} + h^{1/2}(y) (\overline{ds^2})_Y$$

where  $h(y)$  depends only on the coordinates on  $Y$  and satisfies Poisson's equation:

$$\Delta h = \sum \delta(\text{sources}) \quad \text{Kehagias [1998]}$$

- These are generalizations of the solution of [Horowitz, Strominger \[1991\]](#) (for  $Y = \mathbf{C}^3$ ), which leads in the 'near-horizon' limit to the familiar  $AdS_5 \times S^5$  geometry

# Supergravity solutions. III.

- The simplest case is when branes are placed exactly at the singular point
- Then  $(\overline{ds^2})_Y = dr^2 + r^2 \widetilde{ds^2}_X$ , where  $X$  is a Sasaki-Einstein 5-manifold (defined below)
- There is also a 5-form flux through  $X$ :

$$\int_X F_5 \sim N$$

- The function  $h$  is found explicitly:

$$h \sim N \cdot \frac{1}{r^4}, \text{ and the SUGRA ansatz above leads to a smooth geometry of the form}$$

$$AdS_5 \times X^5 \quad \text{Morrison, Plesser [1998]}$$

# $AdS/CFT$ with $\mathcal{N} = 1$ SUSY

- For generic  $X$  such solutions preserve  $N = 1$  SUSY, since  $X$  admits one Killing spinor:

$$(\nabla_\mu + F_5 \gamma_\mu) \psi = 0$$

There are two types of deformations of the above construction:

- One can move the branes off the cone tip
- One can ‘resolve’ the singularity at the tip
- Both of these are related to symmetry breaking in the  $N = 1$  superconformal field theory Klebanov, Witten [1998]

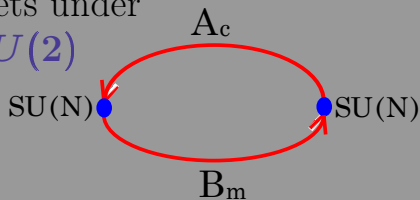
# The conifold theory. I.

- The conifold:  $XY = UV$  in  $\mathbb{C}^4$
- Formal solution:

$$X = a_1 b_1, \quad Y = a_2 b_2, \quad U = a_1 b_2, \quad V = a_2 b_1$$

$\Rightarrow$  Conifold = cone over  $\mathbb{C}P^1 \times \mathbb{C}P^1$

- Dual QFT with gauge group  $SU(N) \times SU(N)$ , two sets of chiral fields:  $A_b$  and  $B_m$ , doublets under global  $SU(2) \times SU(2)$



[Klebanov, Witten \[1998\]](#)

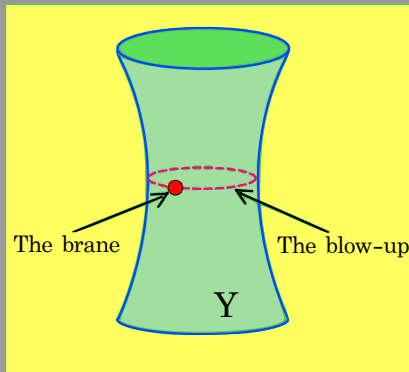


# The conifold theory. II.

Klebanov, Witten [1998]

- The configuration  $[A_c, B_m] = 0$  satisfies the zero-energy condition; diagonalize  $A_c$  and  $B_m$
- If the eigenvalues satisfy  $a_c \neq 0, b_m \neq 0$ , then  $\langle \text{tr}(A_c B_m) \rangle$  can be regarded as positions of the branes moved off the tip of the conifold
- If  $A_c \equiv 0$ , then  $\langle \det B_j \rangle$  may be thought of as positions on the  $\mathbf{CP}^1 \times \mathbf{CP}^1$  glued in at the origin of the resolved cone (hence proportional to the blow-up parameters)

# The conifold theory. III.



# The conifold theory. IV.

- The metric on the resolved conifold (i.e. on  $Y$ ) was built in [Candelas, de la Ossa \[1990\]](#) and generalized in [Pando Zayas, Tseytlin \[2001\]](#)
- A background that interpolates between  $N = 1$  conifold theory (UV) and  $N = 4$  theory after symmetry breaking (IR) (i.e. the  $h$ -function for the brane solution) was constructed in [Klebanov, Murugan \[2007\]](#)

# Geometry of the transverse space $Y$

# Sasaki-Einstein manifolds

- $X$  is Sasaki-Einstein iff the cone over it is Kähler and Ricci-flat:

$$(\overline{ds^2})_Y = dr^2 + r^2 (\widetilde{ds^2})_X$$

- $(\overline{ds^2})_Y$  Kähler & Ricci-flat  $\Leftrightarrow$   
 $(\widetilde{ds^2})_X$  Sasaki-Einstein, of positive curvature

- The metric can be written as

$$(ds^2)_{X^5} = (d\phi - J)^2 + (ds^2)_{\mathcal{M}}$$

where  $(ds^2)_{\mathcal{M}}$  is Kähler-Einstein (but not necessarily smooth),  $J$  is the Kähler current

- $r = 0 \rightarrow$  singularity

# Resolving the singularity of the cone

- It is possible to resolve the singularity of the conical metric by ‘blowing-up’ the vertex, i.e. by replacing it with a cycle of non-zero size
- The metric at infinity, i.e. at  $r \rightarrow \infty$ , will still be asymptotic to the cone:

$$(\overline{ds^2})_Y = dr^2 + r^2 (\widetilde{ds^2})_X \quad \text{for } r \rightarrow \infty$$

- Apart from simplest cases, resolved metrics on the cones are not known  $\Rightarrow$  Our study

# Some examples

- Eguchi, Hanson, 1978

Complex dimension 2, singularity of the form

$$\mathbb{C}^2/Z_2 : (z_1, z_2) \sim (-z_1, -z_2)$$

- Introducing invariant coordinates

$X = z_1^2, Y = z_2^2, Z = z_1 z_2$ , we get an

equation  $XY = Z^2$  in  $\mathbb{C}^3$

- This corresponds to the cone in the embedding of  $\mathbb{CP}^1$  by the linear system  $|\mathcal{O}(2)|$ , i.e. the *anticanonical* embedding

# The Eguchi-Hanson metric

- One can look for the Kähler potential of the form  $K = K(|z_1|^2 + |z_2|^2)$ .  
The metric is, as usual,  $ds^2 = \partial_i \bar{\partial}_j K dz^i d\bar{z}^j$
- For a Kähler metric the Ricci tensor can be expressed as  $R_{i\bar{j}} = -\partial_i \bar{\partial}_j \log \det g$
- Set  $R_{i\bar{j}} = 0$ , solve for the Kähler potential:

## The Eguchi-Hanson metric

$$K = \sqrt{r^2 + 4x^2} + r \log \left( \frac{\sqrt{r^2 + 4x^2} - r}{2x} \right), \quad r > 0$$



# More examples

- 2d case: Eguchi-Hanson = anticanonical cone over  $\mathbf{CP}^1 \Rightarrow \text{SE } X_3 = S^3/\mathbb{Z}_2$
- ‘3d Eguchi-Hanson’ = anticanonical cone over  $\mathbf{CP}^2 \Rightarrow \text{SE } X_5 = S^5/\mathbb{Z}_3$
- 3d case: Candelas-de la Ossa [1990] = anticanonical cone over  $\mathbf{CP}^1 \times \mathbf{CP}^1$  (resolved conifold)  
 $\Rightarrow \text{SE } X_5 := T^{1,1} = \frac{SU(2) \times SU(2)}{U(1)}$

# Other cones?

- One can only build Ricci-flat cones over complex manifolds of ‘positive curvature’ (i.e. with ample anticanonical class)
- For the cone to be of  $\dim_{\mathbb{C}} = 3$ , we take the underlying base to be of  $\dim_{\mathbb{C}} = 2$
- Apart from  $\mathbb{C}P^2$  and  $\mathbb{C}P^1 \times \mathbb{C}P^1$ , there are only 8 other positively curved complex surfaces – the del Pezzo surfaces

$dP_1, \dots, dP_8$

# The del Pezzo surface $dP_1$

- $dP_n$  can be seen as  $CP^2$ , blown-up in  $n$  sufficiently generic points
- We will consider the simplest non-homogeneous case, i.e. the cone over  $dP_1$
- Any metric on  $dP_1$  should have at least two parameters – the sizes of  $CP^2$  and of the blown-up  $CP^1$
- Do these parameters persist in the cone over  $dP_1$ ?

# Isometries

- Whereas the automorphism group of  $\mathbf{CP}^2$  is  $PGL(3, \mathbf{C})$ , the automorphism group of the del Pezzo surface is reduced to

$$Aut(dP_1) = P \begin{pmatrix} \bullet & \bullet & 0 \\ \bullet & \bullet & 0 \\ \bullet & \bullet & \bullet \end{pmatrix} \quad (1)$$

- The isometry group of the metric on the *cone* is the maximal compact subgroup of the parabolic subgroup shown above, i.e.

$$\mathbf{Isom} = U(1) \times U(2)$$

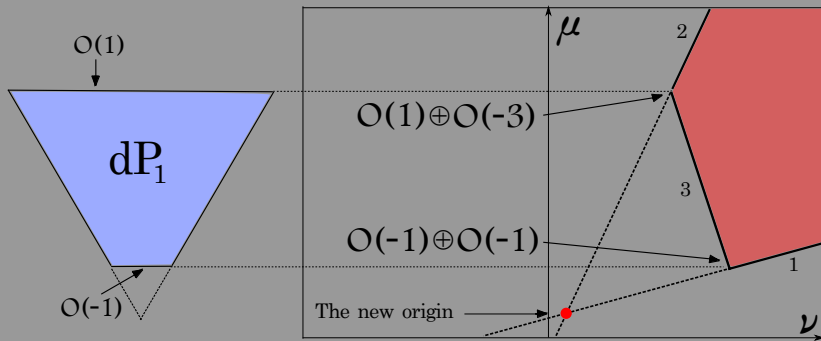
# The main equation

- We will look for a Kähler potential of the form  $K = K(|u|^2, |z_1|^2 + |z_2|^2) := K(e^t, e^s)$
- Just as in the case of the Eguchi-Hanson metric, we can write out a Ricci-flatness equation
- More convenient to perform a Legendre transform w.r.t.  $t, s$ , introducing the dual momentum maps  $\mu = \frac{\partial K}{\partial t}$ ,  $\nu = \frac{\partial K}{\partial s}$  and a dual potential  $G = t\mu + s\nu - K$

# The equation

$$e^{G_\mu + G_\nu} (G_{\mu\mu} G_{\nu\nu} - G_{\mu\nu}^2) = \mu$$

- The domain – the moment polygon



# The expansion at $\infty$

- We can solve the equation exactly at large  $\mu, \nu$  with fixed ‘angle’  $\xi = \frac{\mu}{\nu}$ , assuming the conical form of the metric
- This gives  $G = 3\nu(\log \nu - 1) + \nu P_0(\xi)$
- $P_0(\xi)$  satisfies an ODE and can be found exactly. It provides a Sasaki-Einstein metric, which in the  $dP_1$  case is the  $Y^{2,1}$  manifold ( $Y^{p,q}$  manifolds were constructed in Gauntlett, Martelli, Sparks, Waldram [2004])

# $M^{\text{th}}$ order and the Heun equation

- We can build a systematic perturbation theory

$$G = 3\nu(\log \nu - 1) + \nu P_0(\xi) + \log \nu + \sum_{k=0}^{\infty} \nu^{-k} P_{k+1}(\xi)$$

- In order  $\nu^{-M}$  we obtain the equation

$$\frac{d}{d\xi} \left( Q(\xi) \frac{dP_M}{d\xi} \right) - \left( (M-2)^2 - 1 \right) \xi P_M = \text{r.h.s.},$$

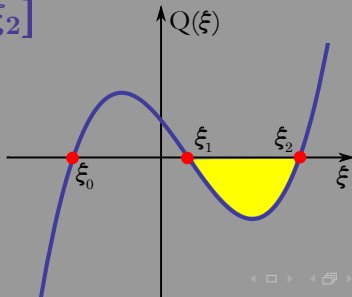
$$\text{where } Q(\xi) = \xi^3 - \frac{3}{2}\xi^2 + d$$

- This is a Heun equation – an analogue of hypergeometric equation with 4 Fuchsian singularities on  $\mathbb{CP}^1$



# Resolution parameters

- All resolution parameters should arise as coefficients in front of the solutions to the homogeneous equation in some order of perturbation theory
- The equation is solved in a ‘physical’ interval  $\xi \in [\xi_1, \xi_2]$

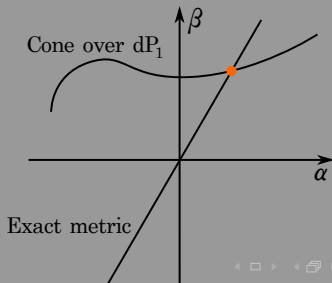


## Resolution parameters. 2.

- Regularity of the metric at the boundaries of the moment polytope requires that the solutions should be regular at  $\xi = \xi_1, \xi_2$   
 $\Rightarrow$  Eigenvalue problem
- Solutions exist for  $M = 3, 4$ :  
 $P_3 = \alpha, P_4 = \beta(\xi - 1)$
- Conjecture:  
For other  $M$  solutions do not exist

# Resolution parameters. 3.

- When  $\beta = -\frac{\alpha}{2\xi_0}$ , the exact metric is known [Calderbank, Gauduchon \[2006\]](#), [Chen, Lu, Pope \[2006\]](#)
- In general, topology imposes one more relation between  $\beta$  and  $\alpha$  [Martelli, Sparks \[2007\]](#)
- Hence the general situation is as follows:



# Questions / Answers

- Can one obtain an exact formula with both parameters  $\alpha, \beta$ ?
- As just discussed, there is an exact formula when  $\beta = -\frac{\alpha}{2\xi_0}$ . Is there a generalization?
- Dual field theories for  $AdS_5 \times X^5$  have been conjectured [Feng, Hanany, He, 2000](#)
- What is the symmetry breaking pattern corresponding to the new parameter in the metric?

Thank you!